

Fri: HW by Spm

Tues: March 28 - HW Spm

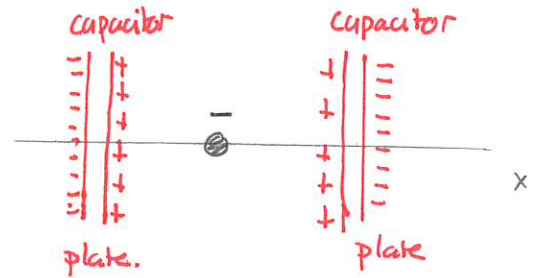
- Read ... Text 5.3, 5.4.

My notes 1.3, 2.0.1.

Position and momentum measurements

We will need to consider particles which can move in one or more spatial dimension. A charged particle trapped in an infinite well could be an example. In this case we could

- * measure the particle's energy
- * measure the " position
- * measure the " momentum
- ...



In order to answer questions about the energy we would need to construct the Hamiltonian. A further question would be

- * how can one construct the Hamiltonian from the comparable classical energy

$$E = \frac{p^2}{2m} + V(x).$$

In order to answer these we will need to describe

- * position states + measurement outcomes
- * position observable
- * momentum states + measurement outcomes
- * momentum observable

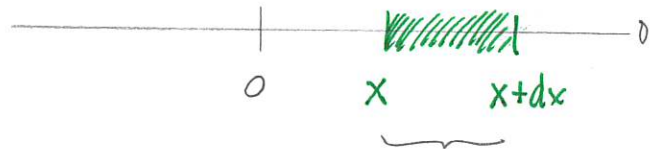
The first complication here is that the measurement outcome possibilities are continuous. This means that we cannot make any statement about the probability of a particular outcome, e.g.

$$\text{Prob}(x = 6.3) = ?$$

$$\text{or Prob}(x = -12.6) = ?$$

Rather we can describe the probability of an outcome in a given range. So we can get

Prob (position is in range $x \rightarrow x+dx$)



what is probability that measurement gives outcome here?

or

Prob ($a \leq x \leq b$)

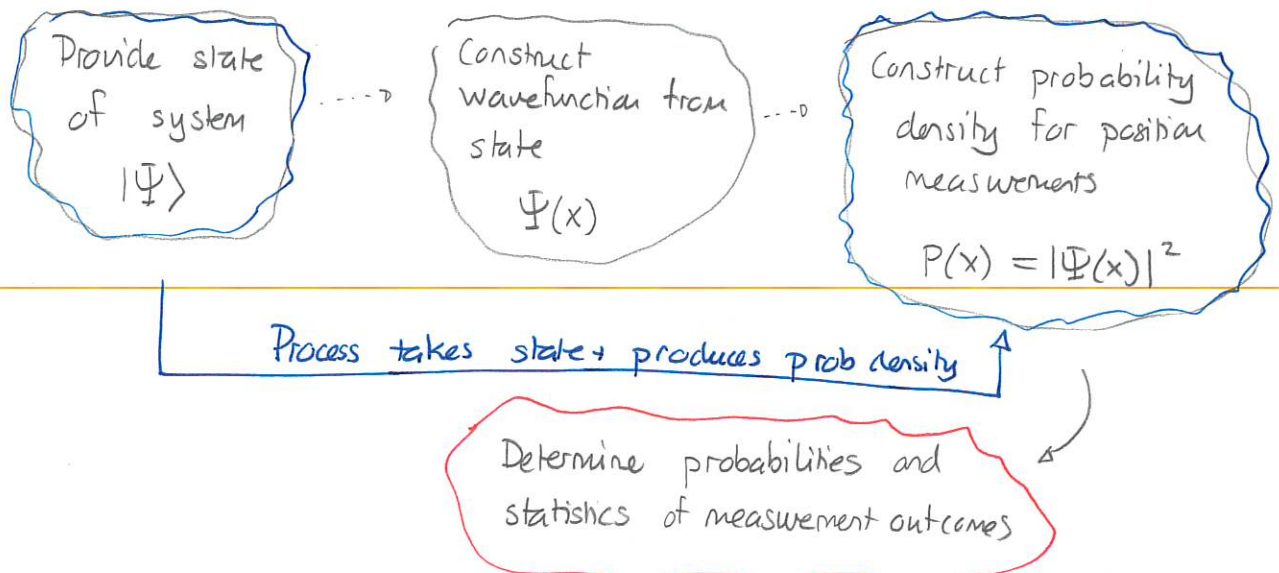
The route to answering these questions is via a probability density $P(x)$. This is a quantity that roughly means

$$\text{Prob (position outcome in range } x \rightarrow x+dx) = P(x) dx$$

or more precisely

$$\text{Prob}(a \leq x \leq b) = \int_a^b P(x) dx$$

What quantum theory will do is.



Position "states" and measurements

Consider a particle restricted to one spatial dimension. In an idealized situation, a position measurement would yield one outcome with perfect precision. We must then have a table

outcome	associated state
\vdots	\vdots
x	$ x\rangle$
\vdots	\vdots

any real number

Then the set of associated states is

$$\{|x\rangle : \text{all real } x\}$$

These must be orthogonal for different outcomes. Thus

$$\boxed{\text{if } x \neq x' \text{ then } \langle x' | x \rangle = 0}$$

We will consider $x = x'$ in a while. These states form a basis and thus we expect that for any state

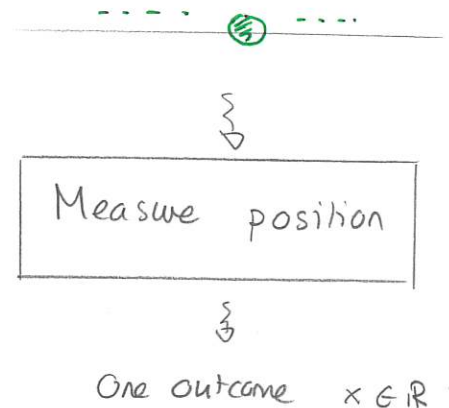
$$|\Psi\rangle = \sum \Psi_x |x\rangle$$

where Ψ_x are real. However the range of x is continuous and the sum must be replaced by an integral. The subscript will be replaced by a functional variable. So

$$\boxed{|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x) |x\rangle dx}$$

sum.
 coefficient basis

where $\Psi(x)$ is a complex valued function (with units m^{-1}) called a wavefunction



Then we obtain a bra via

$$\begin{aligned}\langle \Psi | &= |\Psi\rangle^\dagger \\ &= \left[\int_{-\infty}^{\infty} \Psi(x) |x\rangle dx \right]^\dagger = \int_{-\infty}^{\infty} (\Psi(x) |x\rangle dx)^\dagger \\ &= \int_{-\infty}^{\infty} \Psi^*(x) \langle x| dx.\end{aligned}$$

So

$$\boxed{|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x) |x\rangle dx \quad \Rightarrow \quad \langle \Psi| = \int \Psi^*(x) \langle x| dx}$$

As before the probabilities of measurement outcomes will require inner products. This will require $\langle x'|x\rangle$ for $x' = x$.

Dirac delta function

Recall the Kronecker delta function that has two integer arguments

$$\delta_{mn} = \begin{cases} 1 & m=n \\ 0 & m \neq n. \end{cases}$$

This was useful for certain basis states $\{|n\rangle : n=1,2,\dots\}$ where

$$\langle m|n\rangle = \delta_{mn}$$

This can be extended to continuous variable labels/indices, ... via the

Dirac delta function:

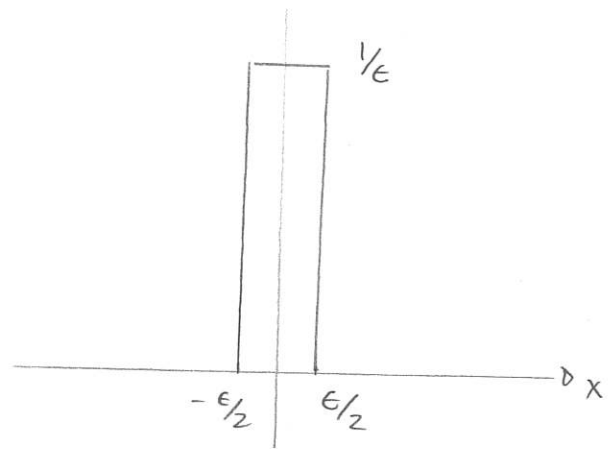
$$\delta(x) = \begin{cases} 0 & \text{if } x \neq 0 \\ \infty & \text{if } x = 0 \end{cases}$$

is defined so that for any function $f(x)$,

$$\int_{-\infty}^{\infty} f(x) \delta(x) dx = f(0)$$

One graphical way to represent uses

$$\delta_\epsilon(x) := \begin{cases} \frac{1}{\epsilon} & -\frac{\epsilon}{2} \leq x \leq \frac{\epsilon}{2} \\ 0 & \text{otherwise} \end{cases}$$



as $\epsilon \rightarrow 0$.

Then definitely

$$\int_{-\infty}^{\infty} \delta_\epsilon(x) dx = \int_{-\infty}^{\infty} \delta_\epsilon(x) 1 dx = 1 = \int_{-\infty}^{\infty} \delta(x) dx$$

and so this at least works for a constant function. For a general function

$$\delta_\epsilon(x) f(x) \rightarrow \frac{1}{\epsilon} f(0) \quad \text{as } \epsilon \rightarrow 0$$

$$\Rightarrow \int_{-\infty}^{\infty} \delta_\epsilon(x) f(x) dx = \int_{-\epsilon/2}^{\epsilon/2} \frac{1}{\epsilon} f(x) dx \rightarrow f(0) \quad \text{as } \epsilon \rightarrow 0$$

1 Dirac delta function: basic rules

a) Show that

$$\delta(x - x') := \begin{cases} 0 & \text{if } x \neq x' \\ \infty & \text{if } x = x'. \end{cases}$$

and

$$\int_{-\infty}^{\infty} f(x') \delta(x - x') dx' = f(x).$$

b) Evaluate

$$\int_{-\infty}^{\infty} x' e^{-x'\lambda} \delta(x - x') dx'.$$

c) Show that

$$\int_{-\infty}^{\infty} \delta(x - x') dx' = 1.$$

Answer: a) If $x \neq x'$ then $\delta(\text{argument} \neq 0) = 0$
if $x = x'$ then $\delta(x - x') = \delta(0) = \infty$

$$\begin{aligned} \int_{-\infty}^{\infty} f(x') \delta(x - x') dx' &= \int_{\infty}^{-\infty} f(x+u) \delta(u) (-du) && u = x - x' \\ &= \int_{-\infty}^{\infty} f(x+u) \delta(u) du = f(x+0) = f(x) \end{aligned}$$

b) Here $f(x') = x' e^{-x'\lambda}$. So

$$\int_{-\infty}^{\infty} x' e^{-x'\lambda} \delta(x - x') dx' = f(x) = x e^{-x\lambda} \Rightarrow \int_{-\infty}^{\infty} x' e^{-x'\lambda} \delta(x - x') dx' =$$

c) Here $f(x') = 1$. So

$$\int_{-\infty}^{\infty} \delta(x - x') dx' = f(x) = 1 \quad \Rightarrow \int_{-\infty}^{\infty} \delta(x - x') dx' = 1$$

With similar derivations:

$$\int_{-\infty}^{\infty} f(x') \delta(x-x') dx' = f(x)$$
$$\int_{-\infty}^{\infty} f(x') \delta(x'-x) dx' = f(x)$$
$$\int_{-\infty}^{\infty} \delta(x'-x) dx = 1$$

Wavefunctions and position states

We now complete the definition of the inner product of position states:

$$\langle x|x' \rangle = \delta(x-x')$$

Then we can extract, formally, the wavefunction from the state.

$$|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x') |x'\rangle dx'$$

$$\Rightarrow \langle x|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x') \langle x|x'\rangle dx' = \int_{-\infty}^{\infty} \Psi(x') \delta(x-x') dx' = \Psi(x)$$

Thus

$$\Psi(x) = \langle x|\Psi\rangle$$

A related useful identity is the completeness relation

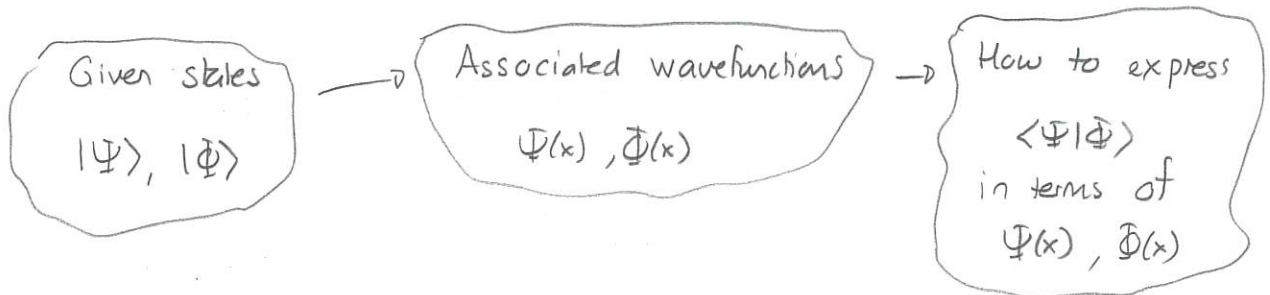
$$\int_{-\infty}^{\infty} |x\rangle\langle x| dx = \hat{I}$$

Proof: $\left[\int_{-\infty}^{\infty} |x\rangle\langle x| dx \right] |\Psi\rangle = \int_{-\infty}^{\infty} |x\rangle \underbrace{\langle x|\Psi\rangle}_{\Psi(x)} dx = \int_{-\infty}^{\infty} \Psi(x) |x\rangle dx = |\Psi\rangle = \hat{I} |\Psi\rangle$ \square

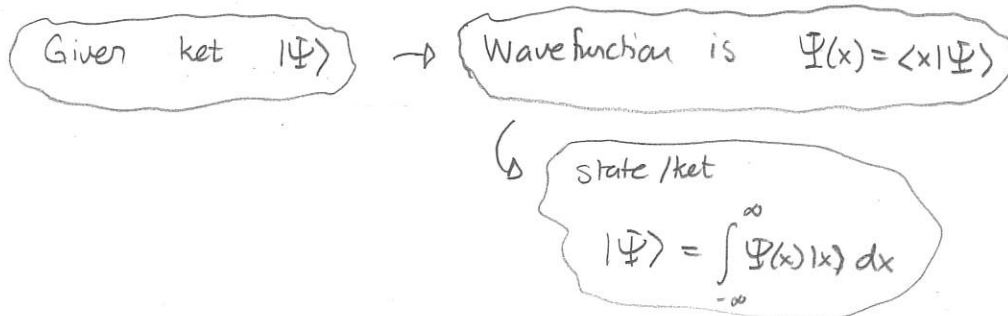
The completeness relation often appears in the midst of calculations and derivations

States, inner products and wavefunctions.

An essential computational requirement is calculating the inner product of two states. We want:



First recall the basic association:



Then

$$|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x)|x\rangle dx$$

$$|\Phi\rangle = \int_{-\infty}^{\infty} \Phi(x)|x\rangle dx \quad \Rightarrow \quad \langle\Phi| = \int_{-\infty}^{\infty} \Phi^*(x)\langle x| dx$$
$$= \int_{-\infty}^{\infty} \Phi^*(x')\langle x'| dx'$$

$$\text{So } \langle\Phi|\Psi\rangle = \int_{-\infty}^{\infty} dx' \Phi^*(x')\langle x'| \int_{-\infty}^{\infty} \Psi(x)|x\rangle dx$$

$$= \int_{-\infty}^{\infty} dx' \int_{-\infty}^{\infty} dx \Phi^*(x') \Psi(x) \underbrace{\langle x'|x\rangle}_{\delta(x'-x)}$$

$$= \int_{-\infty}^{\infty} dx' \Phi^*(x') \underbrace{\int_{-\infty}^{\infty} \Psi(x) \delta(x'-x) dx}_{\Psi(x')}$$

$$= \int_{-\infty}^{\infty} \Phi^*(x') \Psi(x') dx'$$

Thus we have

Given

$$|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x) dx$$
$$|\Phi\rangle = \int_{-\infty}^{\infty} \Phi(x) dx$$

Then

$$\langle \Phi | \Psi \rangle = \int_{-\infty}^{\infty} \Phi^*(x) \Psi(x) dx$$

Note that this gives the normalization requirement

Given

$$|\Psi\rangle = \int_{-\infty}^{\infty} \Psi(x) dx$$

then

$$\langle \Psi | \Psi \rangle = 1 \iff \int_{-\infty}^{\infty} \Psi^*(x) \Psi(x) dx = 1$$
$$\iff \int_{-\infty}^{\infty} |\Psi(x)|^2 dx = 1.$$

Also note that for a bra

$$\langle \Phi | x \rangle = \Phi^*(x)$$

since

$$\langle \Phi | x \rangle = \int_{-\infty}^{\infty} \Phi^*(x') \underbrace{\langle x' | x \rangle}_{\delta(x'-x)} dx'$$
$$= \Phi^*(x)$$

Measurements: Position

If we were able to measure position with perfect precision then we could determine

$$\text{Prob}(\text{position meas} \rightarrow x) = |\langle x | \Psi \rangle|^2 = \langle \Psi | \underbrace{|x\rangle\langle x|}_{\text{position meas. operator}} | \Psi \rangle$$

However we will only be able to do this over a range $a \leq x \leq b$.

So

$$\begin{aligned} \text{Prob}(a \leq x \leq b) &= \langle \Psi | \left[\sum_{a \leq x \leq b} |x\rangle\langle x| \right] | \Psi \rangle \\ &= \langle \Psi | \int_a^b |x\rangle\langle x| dx | \Psi \rangle \\ &= \int_a^b \underbrace{\langle \Psi | x \rangle}_{\Psi^*(x)} \underbrace{\langle x | \Psi \rangle}_{\Psi(x)} dx = \int_a^b \Psi^*(x) \Psi(x) dx. \end{aligned}$$

Thus:

If a system is in state $|\Psi\rangle$ with corresponding wavefunction $\Psi(x)$ then

$$\text{Prob}[\text{position measurement gives } a \leq x \leq b] = \int_a^b |\Psi(x)|^2 dx$$

So the probability density for position measurement outcomes is:

$$P(x) = |\Psi(x)|^2$$

2 Wavefunction mathematics

Consider the states

$$|\Psi_1\rangle \leftrightarrow \Psi_1(x) = Ae^{-|x|/2\alpha^2}$$

$$|\Psi_2\rangle \leftrightarrow \Psi_2(x) = Be^{-|x|/2\beta^2}$$

where $\alpha, \beta > 0$.

a) Determine A and B .

b) Determine $\langle \Psi_1 | \Psi_2 \rangle$.

Answer a) $\langle \Psi_1 | \Psi_1 \rangle = 1$

$$\Rightarrow |A|^2 \int_{-\infty}^{\infty} e^{-|x|/2\alpha^2} e^{-|x|/2\alpha^2} dx = 1$$

$$\Rightarrow |A|^2 \int_{-\infty}^{\infty} e^{-|x|/\alpha^2} dx = 1$$

$$\Rightarrow 2|A|^2 \int_0^{\infty} e^{-x/\alpha^2} dx = 1$$

$$\Rightarrow 2|A|^2 \left[e^{-x/\alpha^2} (-\alpha^2) \Big|_0^{\infty} \right] = 1$$

$$\Rightarrow 2|A|^2 \alpha^2 = 1 \Rightarrow A = \frac{1}{\sqrt{2} \alpha}$$

Similarly $B = \frac{1}{\sqrt{2} \beta}$

$$b) \langle \Psi_1 | \Psi_2 \rangle = \int_{-\infty}^{\infty} A^* e^{-|x|/2\alpha^2} B e^{-|x|/2\beta^2} dx = \frac{2}{2\alpha\beta} \int_0^{\infty} e^{-|x|\gamma} dx$$

$$\text{where } \gamma = \frac{1}{2\alpha^2} + \frac{1}{2\beta^2} = \frac{2\alpha^2 + 2\beta^2}{4\alpha^2\beta^2} = \frac{\alpha^2 + \beta^2}{2\alpha^2\beta^2}$$

$$\Rightarrow \langle \Psi_1 | \Psi_2 \rangle = \frac{1}{\alpha\beta} \frac{1}{\gamma} e^{-|x|\gamma} \Big|_0^{\infty} = \frac{1}{\alpha\beta} \frac{1}{\gamma} = \frac{2\alpha^2\beta^2}{\alpha\beta(\alpha^2 + \beta^2)} = \frac{2\alpha\beta}{\alpha^2 + \beta^2}$$