

Fri HW by 5pm

Tues: Read 9.1, 9.2

HW by 5pm

### Free particles

Consider a particle that can move in one dimension free of any interaction. We can ask

\* what are possible energy eigenstates and eigenvalues for this particle?

mass  $m$



\* are there states such that the particle is localized?

\* how will the state of such a particle evolve with time?

The crucial starting point, that eventually answers questions about energy and time evolution, is to describe the Hamiltonian. Classically the energy is

$$E_{\text{classical}} = \frac{1}{2}mv^2 = \frac{p^2}{2m}.$$

Thus, for a free particle, the Hamiltonian is

$$\hat{H} = \frac{1}{2m} \hat{p}^2 \quad \text{Free particle}$$

where  $\hat{p}$  is the momentum operator. We can then answer questions about energy measurements.

## 1 Free particle energies

Consider a free particle with mass  $m$ .

- Determine the position wavefunction,  $\Psi(x)$ , associated with any energy eigenstate of the particle.
- What are the possible outcomes of an energy measurement?

Answer: a) We need  $|\Psi\rangle$  that satisfied

$$\hat{H}|\Psi\rangle = E|\Psi\rangle$$

$$\Rightarrow \frac{1}{2m} \hat{p}^2 |\Psi\rangle = E|\Psi\rangle$$

Consider a momentum state,  $|p\rangle$ . Then

$$\hat{p}^2 |p\rangle = \hat{p} \hat{p} |p\rangle = \hat{p} p |p\rangle = p \hat{p} |p\rangle = p p |p\rangle = p^2 |p\rangle$$

↑ operator. ↑ number

Thus a possible energy eigenstate is  $|\Psi\rangle = |p\rangle$ . In this case

$$\hat{H}|\Psi\rangle = \frac{p^2}{2m} |\Psi\rangle$$

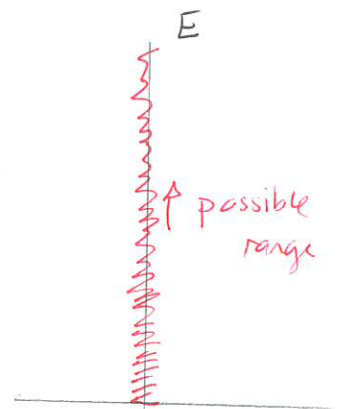
↳ the wavefunction is

$$\Psi(x) = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar}$$

and the energy is:

$$E = \frac{p^2}{2m}$$

- Any real  $p$  is possible. Thus any  $E > 0$  is possible.



So we have that

The energy eigenvalues, eigenstates and wavefunctions for a free particle are indexed by  $p \in \mathbb{R}$ .

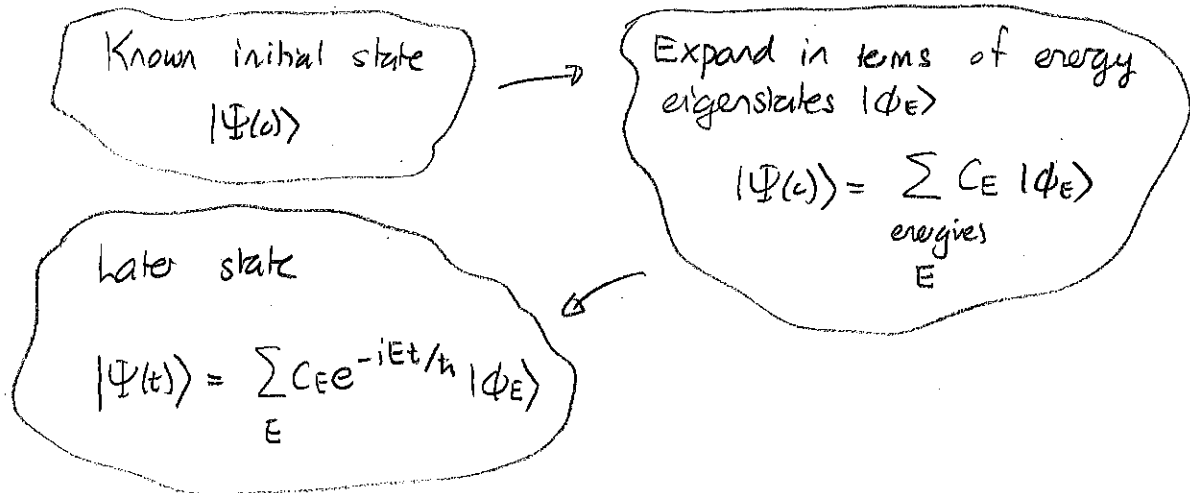
energy	state	wavefunction
$E = \frac{p^2}{2m}$	$ p\rangle$	$\Psi_p(x) = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar}$

### Time evolution for a free particle

Suppose that the initial state and wavefunction for a free particle are known:

$$|\Psi(0)\rangle \rightarrow \Psi(x,t)$$

We would like to be able to determine the state at any later time. We can use a general scheme.



In this case the "summation index" is  $p$  and the summation must be replaced by an integral.

Thus

$$|\Psi(t)\rangle = \int_{-\infty}^{\infty} C(p) |p\rangle dp$$

coefficient ← energy eigenstate

where  $C(p) = \langle p | \Psi(t) \rangle$

$$\Rightarrow C(p) = \tilde{\Psi}(p, 0)$$

$$|\Psi(t)\rangle = \int_{-\infty}^{\infty} C(p) e^{-ip^2 t / 2m\hbar} |p\rangle dp$$

$$= \int_{-\infty}^{\infty} \tilde{\Psi}(p, 0) e^{-ip^2 t / 2m\hbar} |p\rangle dp$$

So

$$\Psi(x, t) = \langle x | \Psi(t) \rangle$$

$$= \int_{-\infty}^{\infty} \tilde{\Psi}(p, 0) e^{-ip^2 t / 2m\hbar} \langle x | p \rangle dp$$

$$\underbrace{e^{ipx/\hbar}}_{\frac{1}{\sqrt{2\pi\hbar}}}$$

Thus the procedure is:

FREE PARTICLE ONLY

1) Determine the initial position space wavefunction

$$\Psi(x, 0)$$

2) Determine the initial momentum space wavefunction

$$\tilde{\Psi}(p, 0) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \Psi(x, 0) dx$$

3) Then at a later time:

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \tilde{\Psi}(p, 0) e^{-ip^2 t / 2m\hbar} e^{ipx/\hbar} dp$$

## 2 Free particle evolution: Gaussian wavefunction

A particle is initially in the state corresponding to the wavefunction

$$\Psi(x, 0) = \left(\frac{1}{\pi a^2}\right)^{1/4} e^{-x^2/2a^2}$$

where  $a > 0$ . This is normalized.

- a) What type of function will represent the momentum wavefunction at  $t = 0$ ? What type of function will represent the momentum wavefunction at any later time? What type of function will represent the position wavefunction at a later time?

The initial momentum wavefunction is

$$\tilde{\Psi}(p, 0) = \left(\frac{a^2}{\pi \hbar^2}\right)^{1/4} e^{-p^2 a^2 / 2\hbar^2}$$

- b) Determine the position wavefunction and position probability density at any later time.

Note the following integrals, all true if the real part of  $\alpha$  is positive.

$$\begin{aligned} \int_{-\infty}^{\infty} e^{-\alpha x^2 + \beta x + \gamma} dx &= e^{\gamma} e^{\beta^2/4\alpha} \sqrt{\frac{\pi}{\alpha}} \\ \int_{-\infty}^{\infty} x e^{-\alpha x^2 + \beta x + \gamma} dx &= e^{\gamma} e^{\beta^2/4\alpha} \frac{\beta}{2} \sqrt{\frac{\pi}{\alpha^3}} \\ \int_{-\infty}^{\infty} x^2 e^{-\alpha x^2 + \beta x + \gamma} dx &= e^{\gamma} e^{\beta^2/4\alpha} \frac{\beta^2 + 2\alpha}{4} \sqrt{\frac{\pi}{\alpha^5}} \end{aligned}$$

Answer: a) 
$$\begin{aligned} \tilde{\Psi}(p, 0) &= \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \Psi(x, 0) e^{-ipx/\hbar} dx \\ &= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{1}{\pi a^2}\right)^{1/4} \int e^{-x^2/2a^2} e^{-ipx/\hbar} dx \end{aligned}$$

This is a Gaussian integral. The result will be another Gaussian function of  $p$ . At a later time

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \tilde{\Psi}(p, 0) e^{ipx/\hbar} e^{-ip^2 t/2m\hbar} dp$$

Gaussian

will also be Gaussian.

$$\begin{aligned}
 \text{b) } \Psi(x,t) &= \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} \tilde{\Psi}(p,0) e^{ipx/\hbar} e^{-ip^2t/2m\hbar} dp \\
 &= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{a^2}{\pi\hbar^2}\right)^{1/4} \int_{-\infty}^{\infty} e^{-p^2 a^2/2\hbar^2} e^{-ip^2t/2m\hbar} e^{ipx/\hbar} dp \\
 &= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{a^2}{\pi\hbar^2}\right)^{1/4} \int_{-\infty}^{\infty} e^{-p^2 \left(\frac{a^2}{2\hbar^2} + \frac{it}{2m\hbar}\right) + p \left(\frac{ix}{\hbar}\right)} dp
 \end{aligned}$$

The integral has form

$$\int_{-\infty}^{\infty} e^{-p^2\alpha + p\beta + \gamma} dp$$

$$\text{where } \alpha = \left(\frac{a^2}{2\hbar^2} + \frac{it}{2m\hbar}\right) = \frac{ma^2 + it\hbar}{2m\hbar^2}$$

$$\beta = \frac{ix}{\hbar}$$

$$\gamma = 0$$

Since  $\text{Re}(\alpha) > 0$  the formulas apply. Then

$$\int_{-\infty}^{\infty} e^{-p^2\alpha + p\beta + \gamma} dp = e^{\beta^2/4\alpha} \sqrt{\frac{\pi}{\alpha}}$$

$$\text{So } \frac{\beta^2}{4\alpha} = \frac{-x^2}{4\hbar^2} \frac{(2m\hbar^2)}{(ma^2 + it\hbar)} = \frac{-mx^2}{2} \frac{1}{ma^2 + it\hbar}$$

$$\text{But } \frac{1}{ma^2 + it\hbar} = \frac{ma^2 - it\hbar}{ma^2 - it\hbar} \frac{1}{ma^2 + it\hbar} = \frac{ma^2 - it\hbar}{m^2a^4 + t^2\hbar^2}$$

So

$$\frac{\beta^2}{4\alpha} = -\frac{m^2 x^2}{2} \quad \frac{m^2 a^2 - i\hbar t}{m^2 a^2 + i\hbar^2 t} = \frac{-i\hbar m x^2 t}{2(\dots)} - \frac{m^2 a^2}{2(m^2 a^2 + i\hbar^2 t)} x^2$$

Thus:

$$\Psi(x,t) = \frac{1}{\sqrt{2\pi\hbar^2}} \left(\frac{a^2}{\pi\hbar^2}\right)^{1/4} \underbrace{e^{-i(\dots)}}_{\text{Overall phase}} e^{-x^2 \frac{m^2 a^2}{2(m^2 a^2 + i\hbar^2 t)}} \sqrt{\frac{2\pi\hbar^2 m}{m^2 a^2 + i\hbar^2 t}}$$

$$= \left(\frac{a^2}{\pi}\right)^{1/4} \sqrt{\frac{m}{m^2 a^2 + i\hbar^2 t}} e^{-i(\dots)} e^{-x^2 \frac{m^2 a^2}{2(m^2 a^2 + i\hbar^2 t)}}$$

$$= \left(\frac{a^2}{\pi}\right)^{1/4} \underbrace{\sqrt{\frac{m}{m^2 a^2 + i\hbar^2 t}}}_{\text{phase}} e^{-i(\dots)} \underbrace{e^{-x^2 / 2a^2 \left(1 + \frac{i\hbar^2 t}{m^2 a^2}\right)}}_{\text{Gaussian.}}$$

Then the position probability density is:

$$P(x,t) = |\Psi(x,t)|^2 = \sqrt{\frac{a^2}{\pi}} \frac{m}{\sqrt{(m^2 a^2 + i\hbar^2 t)(m^2 a^2 - i\hbar^2 t)}} e^{-x^2 / a^2 \left(1 + \frac{i\hbar^2 t}{m^2 a^2}\right)}$$

This describes a Gaussian with width

$$a \sqrt{1 + \frac{i\hbar^2 t}{m^2 a^2}}$$

We see that as time passes the width of the distribution increases

Demo: GuV's Gaussian Wave Packet

Note that for the free particle

$$\tilde{\Psi}(p, 0) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx'/\hbar} \Psi(x', 0) dx'$$

gives:

$$\Psi(x, t) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dp \int_{-\infty}^{\infty} dx' \Psi(x', 0) e^{-ipx'/\hbar} e^{ipx/\hbar} e^{-ip^2 t/2m\hbar}$$

$$\Psi(x, t) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dx' \int_{-\infty}^{\infty} dp e^{ip(x-x')/\hbar} e^{-ip^2 t/2m\hbar} \Psi(x', 0)$$

later  
wavefunction

\* can be integrated

\* function of  $x, x'$

~~initial~~  
initial  
wavefunction

amounts to evolution  
operator