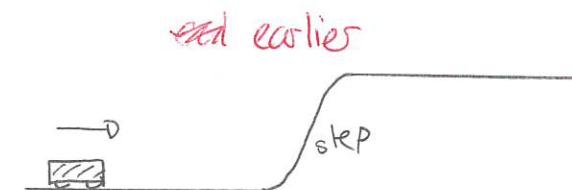


Weds 6.1, 6.2

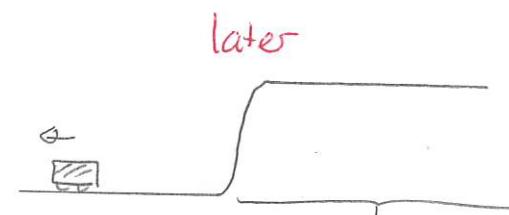
Particles at step and barrier potentials

So far we have applied quantum theory to completely free particles and also bound particles. We now consider situations where particles encounter steps or barriers. These will demonstrate unusual properties of quantum theory that are absent in classical physics. We can illustrate the classical via a cart sliding along a track.

First consider a step. Here there are two possibilities



moves freely with  
slower speed

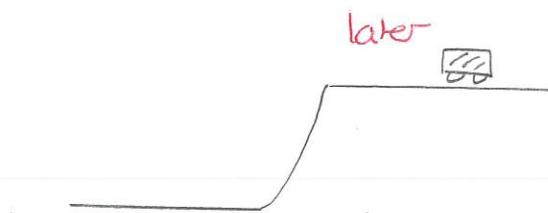


later **definitely**  
returns with  
**certainty**  
moves freely

OR



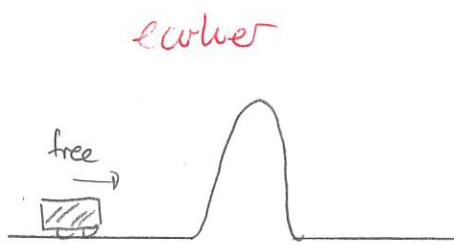
moves freely with a  
faster speed



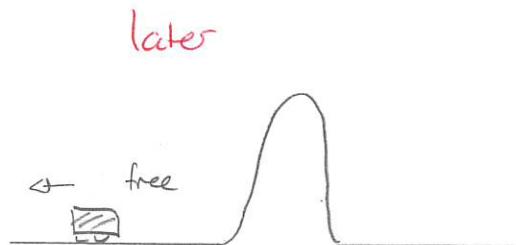
**never** returns  
to this region

**definitely** enters this  
region, moves freely

Second consider a barrier



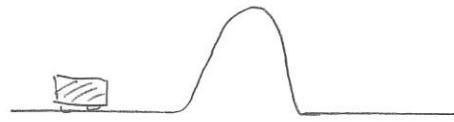
moves freely with  
slower speed



later definitely  
returns with  
certainty

never enters this  
region

OR



moves freely  
with higher  
speed



never  
returns.

later moves  
freely  
definitely enters this  
region

We shall analyse these situations in quantum theory and see that the "never," "definitely" and "certainty" terms no longer apply.

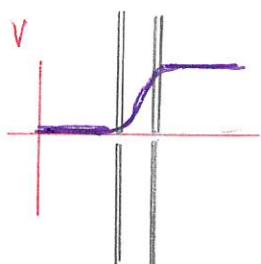
The procedure will be to:

- 1) provide a potential that describes the step /barrier
- 2) solve the TISE for the potential
- 3) interpret the solutions in terms of "transmission" and "reflection."

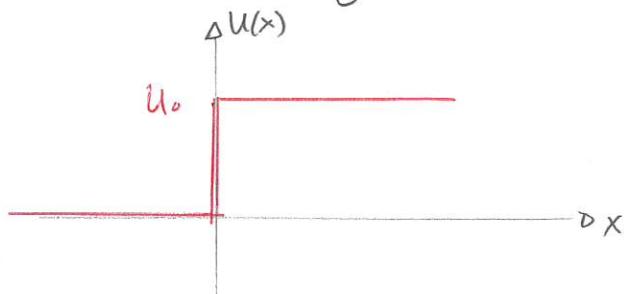
The barrier potential particularly has many technical and fundamental applications.

We could imagine creating such potentials with various electrodes that are charged. For example a step potential might involve. Mathematically this

may be describe in terms of potentials that are piecewise constant with sudden jumps.



0V 10V



### Free particle solutions

In classical physics in any region where the potential is constant the particle is free and moves with constant momentum. In one dimension this can either be left or right. This is described by the sign of the momentum.

In quantum theory the simplest such case occurs when  $U(x)=0$ . Then the TISE gives

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi(x)$$

The possible solutions to this are

$$\psi(x) = Ae^{ikx} \quad \text{or} \quad \psi(x) = Ae^{-ikx}$$

where we require  $k \geq 0$

**Quiz 1**  $\rightarrow$  Type-particle

Substitution gives

$$-\frac{\hbar^2}{2m} (\pm k)^2 = E \Rightarrow \frac{\hbar^2 k^2}{2m} = E \Rightarrow k = \sqrt{\frac{2mE}{\hbar^2}}$$

$$\text{But } E = \frac{p^2}{2m} \Rightarrow p^2 = \hbar^2 k^2 \Rightarrow |p| = \hbar k$$

We need to resolve the sign of the momentum. We will have to stipulate how this is implemented. By convention we will always operate with  $k \geq 0$ .

Consider the momentum operator acting on each. Thus

$$\hat{p}\Psi = -i\hbar \frac{\partial \Psi}{\partial x}$$

If  $\Psi = Ae^{ikx}$  then this gives

$$\hat{p}\Psi = -i\hbar(ik)\Psi = \hbar k\Psi \rightarrow \text{suggests momentum is positive } \hbar k$$

If  $\Psi = Ae^{-ikx}$  then this gives

$$\hat{p}\Psi = -i\hbar(-ik)\Psi = -\hbar k\Psi \rightarrow \text{suggests momentum is negative: } -\hbar k.$$

Thus

Free particles with energy  $E$  are described by (here  $k = \sqrt{\frac{2mE}{\hbar^2}}$ )

$$\Psi(x) = Ae^{ikx}$$

$$\text{momentum } p = \hbar k$$

travels right



$$\Psi(x) = Ae^{-ikx}$$

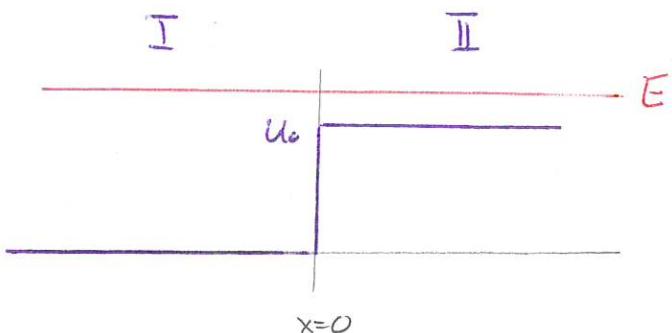
$$\text{momentum } p = -\hbar k$$

travels left

## Quantum theory for a step potential (energy greater than step)

A step potential has the form

$$U(x) = \begin{cases} 0 & x < 0 \\ U_0 & x > 0 \end{cases}$$

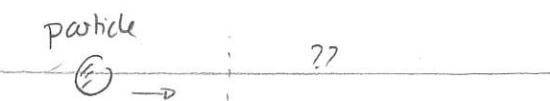


We consider the situation where

the total energy  $E > U_0$ . We

assume that the particle is

incident from the left. In classical physics we expect that it will definitely move right beyond the step but with reduced momentum.



In quantum theory we solve the TISE in each region:

$$\text{Region I } (x < 0) \quad -\frac{\hbar^2}{2m} \frac{d^2\psi_I}{dx^2} = E\psi_I$$

$$\text{Region II } (x > 0) \quad -\frac{\hbar^2}{2m} \frac{d^2\psi_{II}}{dx^2} + U_0\psi_{II} = E\psi_{II}$$

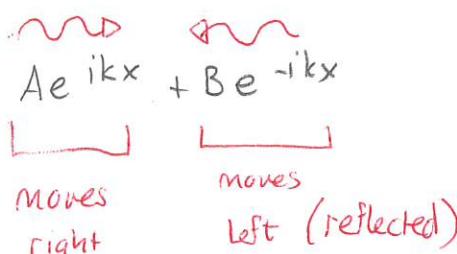
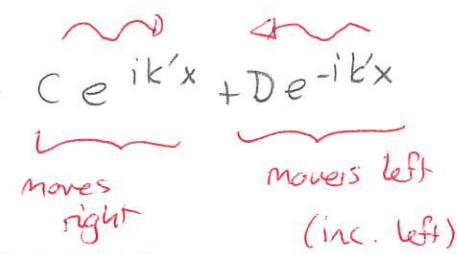
$$\Rightarrow -\frac{\hbar^2}{2m} \frac{d^2\psi_{II}}{dx^2} = (E - U_0)\psi_{II}$$

We will also require that, at the boundary

$$1) \quad \psi_I(0) = \psi_{II}(0) \quad (\text{wavefunctions match})$$

$$2) \quad \left. \frac{d\psi_I}{dx} \right|_0 = \left. \frac{d\psi_{II}}{dx} \right|_0 \quad (\text{slopes match})$$

Then the generic solutions are:

Region I	Region II
$\Psi_I = Ae^{ikx} + Be^{-ikx}$  (incident) $k = \sqrt{\frac{2mE}{\hbar^2}}$	$\Psi_{II} = Ce^{ik'x} + De^{-ik'x}$  transmitted $k' = \sqrt{\frac{2m(E-U_0)}{\hbar^2}}$

We assume that particles are incident from the left. Thus  $D=0$ . So

$$\begin{aligned}\Psi_I &= Ae^{ikx} + Be^{-ikx} & k &= \sqrt{2mE/\hbar^2} \\ \Psi_{II} &= Ce^{+ik'x} & k' &= \sqrt{2m(E-U_0)/\hbar^2}\end{aligned}$$

Now we match at the boundary.

Quiz  $\rightarrow$  change  $-ik'$  to  $ik'$

Thus

$$\begin{aligned}A + B &= C \\ A - B &= \frac{k'}{k} C\end{aligned}$$

linear algebra gives:

$$\boxed{\begin{aligned}B &= \frac{k-k'}{k+k'} A \\ C &= \frac{2k}{k+k'} A\end{aligned}}$$

$\rightarrow$  some particles will be reflected back

$\rightarrow$  some particles (not all) will be transmitted.