

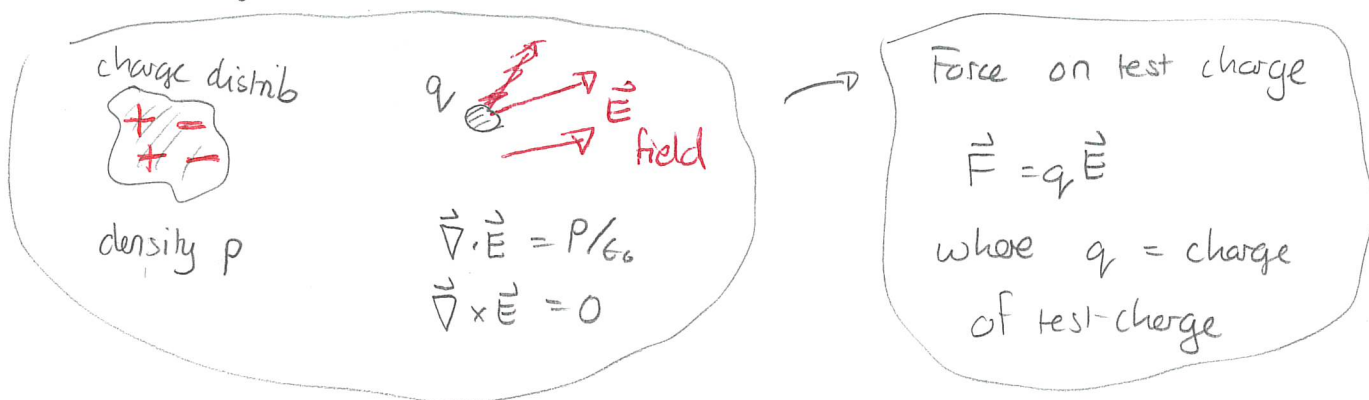
Fri: HW 16

Mon: Read 5.1.3

Tues:

Magnetostatics

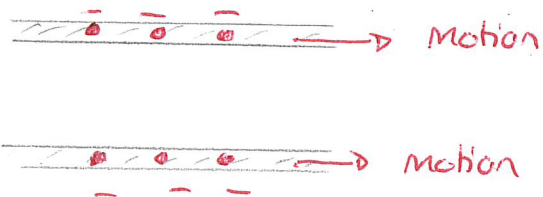
Electrostatics describes the forces that stationary source charges exert on test charges. It does so via electric fields.



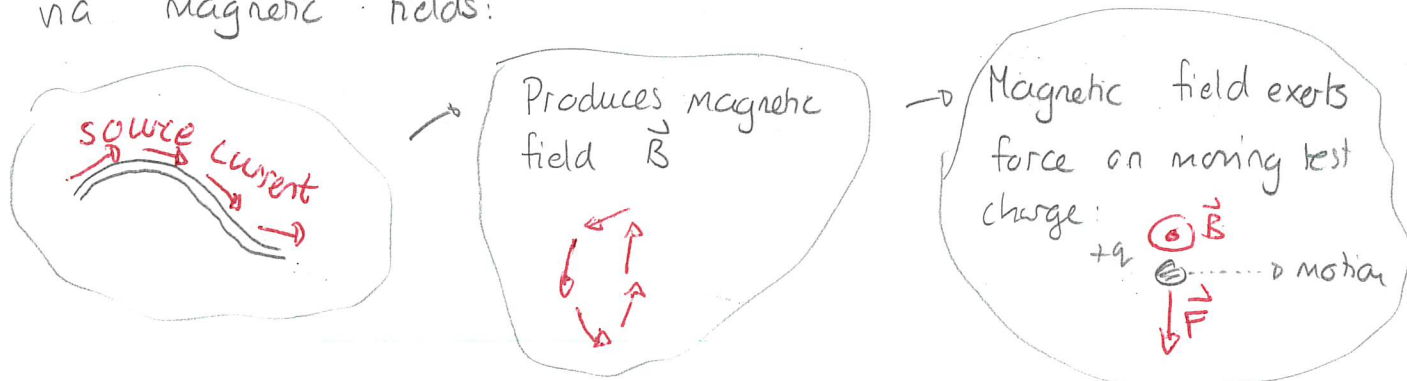
This does not describe forces exerted by moving charges. Consider two parallel identical currents. We have

\* Coulomb's law predicts  $\Rightarrow$  repel

\* Observation  $\Rightarrow$  attract.



There is clearly a different process here. We will describe this via magnetic fields:



There are two parts to this:

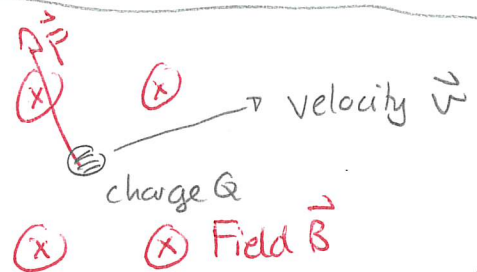
- 1) describe how a magnetic field exerts a force on a moving test charge
- 2) describe how currents produce magnetic fields.

### Force exerted by a magnetic field

We will assume that there is a mechanism for describing how magnetic fields are produced. An indirect definition involves the force exerted on a moving charged particle. Observations eventually lead to:

If a particle with charge  $Q$  moves with velocity  $\vec{v}$  through a magnetic field  $\vec{B}$  then the force exerted by the field on the charge is

$$\vec{F} = Q \vec{v} \times \vec{B}$$



Note that:

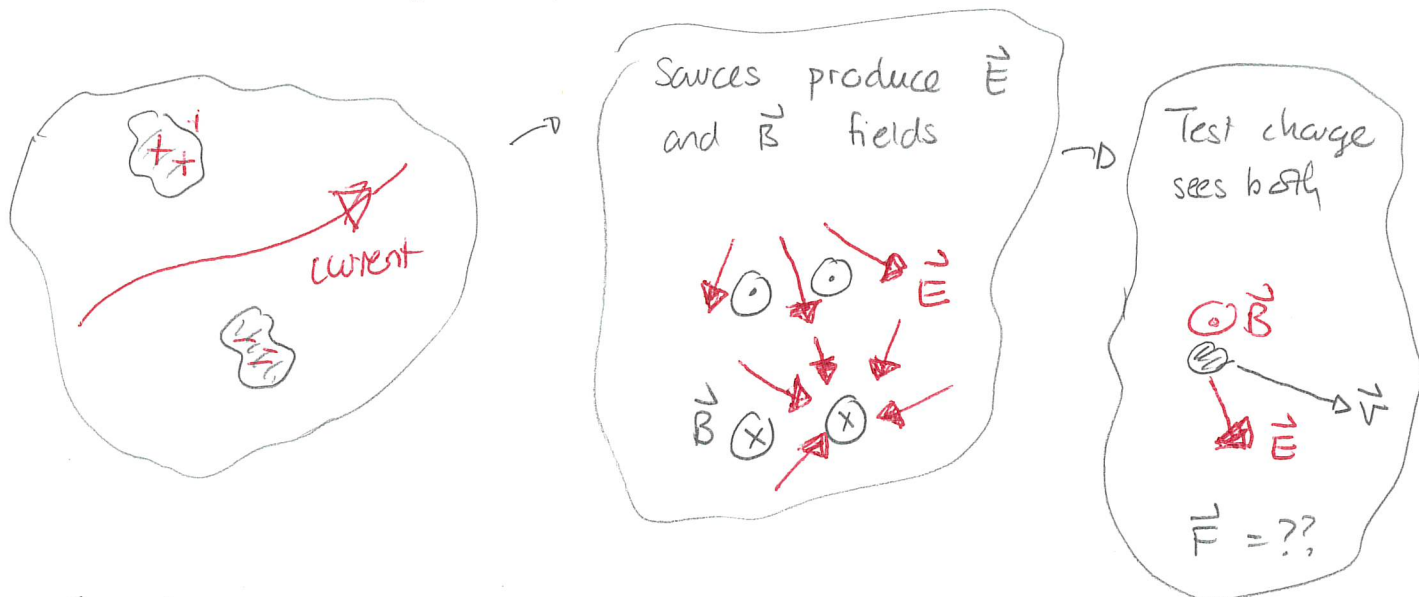
- 1) the force produced by a magnetic field is perpendicular to the field
- 2) the force tends to produce changes in the direction of motion
- 3) the units satisfy

$$N = C \text{ m/s} \times \text{units of } B$$

$$\Rightarrow \text{units of } B = \frac{Ns}{C \cdot m} = \frac{N}{A \cdot m} \equiv \text{Tesla} \equiv T$$

## Lorentz Force Law

Suppose that a test particle is subjected to both electric and magnetic fields. These each exert a force on the particle and the two forces can be combined.



The fundamental rule that describes this is the Lorentz Force Law

Suppose a test particle with charge  $Q$  is in an electric and magnetic field. Then the force exerted on the test charge is

$$\vec{F} = Q(\vec{E} + \vec{v} \times \vec{B})$$

where  $\vec{v}$  = velocity of particle

$\vec{E}$  = electric field at particle location

$\vec{B}$  = magnetic " " " "

This is the fundamental force law in electromagnetic theory. Then Newton's 2<sup>nd</sup> Law,  $\vec{F}_{\text{net}} = m\vec{a}$  determines acceleration

## 1 Particle in a uniform magnetic field

A particle with mass  $m$  and charge  $Q$  is placed in a uniform magnetic field  $\mathbf{B} = B\hat{z}$  (here  $B$  is independent of location and time). The electric field in this region is  $\mathbf{E} = 0$ . We aim to convert the equations of motion for this particle into a set of differential equations and to then solve these, knowing the initial velocity.

- a) Let  $\mathbf{v} = v_x(t)\hat{x} + v_y(t)\hat{y} + v_z(t)\hat{z}$  be the velocity of the particle. Combine Newton's second law and the Lorentz force law to produce a set of coupled differential equations for  $v_x(t)$ ,  $v_y(t)$  and  $v_z(t)$ . Rewrite these in terms of the cyclotron frequency

$$\omega := \frac{QB}{m}.$$

- b) Solve the differential equation for  $v_z(t)$  in terms of the initial velocity  $v_z(0)$ .
- c) Manipulate the two first order differential equations for  $v_x$  and  $v_y$  to produce uncoupled second order differential equations for each of these.
- d) Solve each of these to yield  $v_x(t)$  and  $v_y(t)$  in terms of  $v_x(0)$  and  $v_y(0)$ .
- e) Use these to find expressions for the components of the position of the particle,  $x(t)$ ,  $y(t)$  and  $z(t)$  (in terms of initial positions and velocities).
- f) Describe the trajectory followed by the particle geometrically. What does the cyclotron frequency describe in terms of the particle motion?

Answer: a)  $\vec{F}_{\text{net}} = m\vec{a}$

$$\Rightarrow m \frac{d\vec{v}}{dt} = \vec{F}_{\text{net}}$$

$$\Rightarrow m \frac{d\vec{v}}{dt} = Q \vec{v} \times \vec{B}$$

Now  $\vec{v} = v_x \hat{x} + v_y \hat{y} + v_z \hat{z}$

$$\vec{B} = B \hat{z}$$

$$\Rightarrow \vec{v} \times \vec{B} = v_x B \underbrace{\hat{x} \times \hat{z}}_{-\hat{y}} + v_y B \underbrace{\hat{y} \times \hat{z}}_{\hat{x}} + v_z B \hat{z} \times \hat{z} = 0$$

So  $Q \vec{v} \times \vec{B} = Q v_y B \hat{x} - Q v_x B \hat{y}$

But  $\frac{d\vec{v}}{dt} = \frac{dv_x}{dt} \hat{x} + \frac{dv_y}{dt} \hat{y} + \frac{dv_z}{dt} \hat{z}$

$$\begin{aligned} \text{Thus } m \frac{d\vec{v}}{dt} &= m \frac{dv_x}{dt} \hat{x} + m \frac{dv_y}{dt} \hat{y} + m \frac{dv_z}{dt} \hat{z} \\ &= Qv_y B \hat{x} - Qv_x B \hat{y} \end{aligned}$$

Comparing coefficients gives:

$$\left. \begin{aligned} m \frac{dv_x}{dt} &= QB v_y \\ m \frac{dv_y}{dt} &= -QB v_x \\ m \frac{dv_z}{dt} &= 0 \end{aligned} \right\} \Rightarrow \begin{cases} \frac{dv_x}{dt} = \omega v_y \\ \frac{dv_y}{dt} = -\omega v_x \\ \frac{dv_z}{dt} = 0 \end{cases}$$

b)  $v_z(t) = \text{constant} \Rightarrow v_z(t) = v_z(0) \equiv \text{constant}$

c)  $\frac{d}{dt} \frac{dv_x}{dt} = \omega \frac{dv_y}{dt}$

$$\Rightarrow \frac{d^2 v_x}{dt^2} = \omega (-v_x) \Rightarrow \frac{d^2 v_x}{dt^2} = -\omega^2 v_x$$

$$\frac{d}{dt} \left( \frac{dv_y}{dt} \right) = -\omega \frac{dv_x}{dt} \Rightarrow \frac{d^2 v_y}{dt^2} = -\omega^2 v_y$$

d) Consider

$$\frac{d^2 v_x}{dt^2} = -\omega^2 v_x$$

The general solution is

$$v_x(t) = A \cos \omega t + B \sin \omega t$$

where A B are constants

We could attempt to construct a similar solution for  $v_y$  independently.  
But  $v_x, v_y$  are related. For example

$$\frac{dv_x}{dt} = \omega v_y \Rightarrow v_y = \frac{1}{\omega} \frac{dv_x}{dt}$$

$$\begin{aligned} \text{Thus } v_y(t) &= \frac{1}{\omega} [-A\omega \sin \omega t + B\omega \cos \omega t] \\ &= -A \sin \omega t + B \cos \omega t \end{aligned}$$

So we have a general solution:

$$v_x(t) = A \cos \omega t + B \sin \omega t$$

$$v_y(t) = -A \sin \omega t + B \cos \omega t$$

$$\text{Then } v_x(0) = A \underbrace{\cos 0}_1 + B \underbrace{\sin 0}_0 \Rightarrow A = v_x(0)$$

$$v_y(0) = -A \sin 0 + B \underbrace{\cos 0}_1 \Rightarrow B = v_y(0)$$

Thus the solution is:

$$v_x(t) = v_x(0) \cos(\omega t) + v_y(0) \sin(\omega t)$$

$$v_y(t) = -v_x(0) \sin(\omega t) + v_y(0) \cos(\omega t)$$

$$v_z(t) = v_z(0)$$

e) We have

$$v_x(t) = \frac{dx}{dt} = v_x(0) \cos(\omega t) + v_y(0) \sin(\omega t)$$

$$v_y(t) = \frac{dy}{dt} = -v_x(0) \sin(\omega t) + v_y(0) \cos(\omega t)$$

$$v_z(t) = \frac{dz}{dt} = v_z(0)$$

Each can be integrated independently. So

$$x(t) = \frac{v_x(0)}{\omega} \sin(\omega t) - \frac{v_y(0)}{\omega} \cos(\omega t) + C_x \leftarrow \text{constant}$$

$$y(t) = \frac{v_x(0)}{\omega} \cos(\omega t) + \frac{v_y(0)}{\omega} \sin(\omega t) + C_y$$

$$z(t) = v_z(0)t + C_z$$

The remaining constants can be fixed in terms of initial locations So

$$x(0) = \frac{v_x(0)}{\omega} \sin(0) - \frac{v_y(0)}{\omega} \cos(0) + C_x$$

$$\Rightarrow C_x = x(0) + \frac{v_y(0)}{\omega}$$

$$y(0) = \frac{v_x(0)}{\omega} \cos(0) + \frac{v_y(0)}{\omega} \sin(0) + C_y$$

$$\Rightarrow C_y = y(0) - \frac{v_x(0)}{\omega}$$

$$z(0) = v_z(0) \cdot 0 + C_z$$

$$\Rightarrow C_z = z(0)$$

We can rewrite these in terms of

$$V_x(0) = V_{0x}$$

$$x(0) = x_0$$

etc...

Thus

$$x(t) = \frac{V_{0x}}{\omega} \sin(\omega t) - \frac{V_{0y}}{\omega} \cos(\omega t) + \frac{V_{0y}}{\omega} + x_0$$

$$y(t) = \frac{V_{0x}}{\omega} \cos(\omega t) + \frac{V_{0y}}{\omega} \sin(\omega t) - \frac{V_{0x}}{\omega} + y_0$$

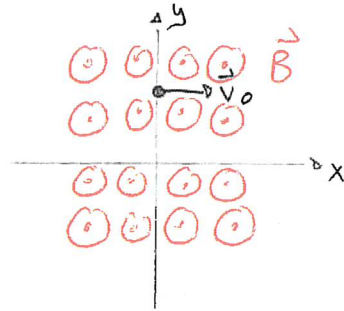
$$z(t) = V_{0z} t + z_0$$

f) Consider the simplest case:

$$V_{0y} = 0 \quad z_0 = 0$$

$$V_{0z} = 0$$

$$V_{0x} = V_0 \neq 0$$



Then:

$$x(t) = \frac{V_0}{\omega} \sin(\omega t) + x_0$$

$$y(t) = \frac{V_0}{\omega} \cos(\omega t) - \frac{V_0}{\omega} + y_0$$

$$z(t) = 0$$

This is an orbit in the  $xy$  plane since  $z=0$



Then

$$x(t) - x_0 = \frac{V_0}{\omega} \sin(\omega t)$$

$$y(t) - y_0 + \frac{V_0}{\omega} = \frac{V_0}{\omega} \cos(\omega t)$$

let  $\alpha = x_0$

$$\beta = y_0 - \frac{V_0}{\omega}$$

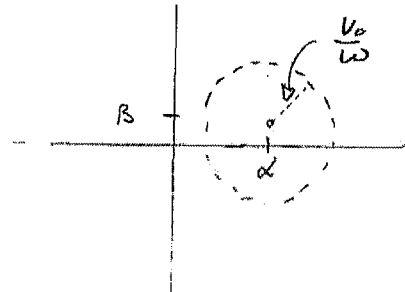
Then  $x(t) - \alpha = \frac{V_0}{\omega} \sin(\omega t)$

$$y(t) - \beta = \frac{V_0}{\omega} \cos(\omega t)$$

So  $(x(t) - \alpha)^2 + (y(t) - \beta)^2 = \left(\frac{V_0}{\omega}\right)^2$ . This describes a circle with radius  $\frac{V_0}{\omega}$  centered at  $(\alpha, \beta)$ .

The orbit is counterclockwise with angular frequency

$$\omega = \frac{QB}{m}$$



Such circular motion is common in uniform magnetic fields.

The motion is more complicated if there is a component of velocity along the field direction. In this case it is helical

Demo: PSU - S

Charge in Unif field I

Charge in unif field II

Charge in bottle etc...