

HW 5 due Monday

Magnetic Fields in Electromagnetic Waves

Recall that electromagnetic fields in a vacuum satisfy

$$\begin{aligned} \vec{\nabla} \cdot \vec{E} &= 0 & \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} & \vec{\nabla} \times \vec{B} &= \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \end{aligned}$$

These can be combined to yield wave equations:

$$\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\nabla^2 \vec{B} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

Consider harmonic plane wave solutions, which, for the electric field have the form

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t + \phi)}$$

This has the properties:

i) the angular frequency + wavenumber are related by:

$$c = \omega/k \quad \text{where} \quad c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

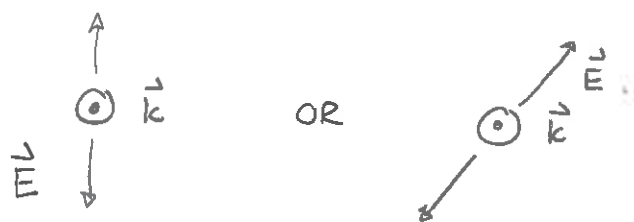
is the speed with which these waves propagate.

- ii) the direction along which waves propagate is \vec{k}
- iii) the equation $\vec{\nabla} \cdot \vec{E}$ implies that $\vec{k} \cdot \vec{E} = 0$. Thus

At every location, the electric field is perpendicular to the direction of propagation.

This is thus a transverse wave.

- iv) If \vec{E}_0 is independent of \vec{r} and t then at every location the electric field lies along the same line. At any single location it oscillates back + forth along one line. It is the



called linearly polarized and the direction of polarization is determined by \vec{E}_0 .

Demo: PSU Schuylkill demos

The magnetic field satisfies a similar equation and we might expect that harmonic wave solutions exist. This is true. However, it is important to note that the magnetic field depends on the electric field since

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Exercise 1 Consider an electromagnetic wave that has the following properties:

- i) It propagates along the \hat{z} direction with wavelength 3.0m
- ii) It is linearly polarized along the x direction, and the electric field has amplitude 20 N/C .

The purpose of this exercise is to show that these completely specify the electric and magnetic fields for the wave.

a) The electric field has form

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

Determine expressions for \vec{k} , ω and \vec{E}_0 .

b) Using $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$, determine an expression for the magnetic field. Ensure that you find expression for its frequency, wave vector, and the amplitude (a vector).

c) Show that

$$\omega \vec{B} = \vec{k} \times \vec{E}$$

Answer: i) First $\vec{k} = k_x \hat{i} + k_y \hat{j} + k_z \hat{k}$ gives the direction of propagation. Thus $k_x = k_y = 0$ and

$$\vec{k} = k \hat{k} \quad \text{with vector along } \hat{z}$$

Now $k = 2\pi/\lambda = \frac{2\pi}{3.0} \text{ m}^{-1} = 2.1 \text{ m}^{-1}$

So $\vec{k} = 2.1 \text{ m}^{-1} \hat{k}$

Then $v = \omega/k \Rightarrow \omega = vk$

$$= 3.0 \times 10^8 \text{ m/s} \times \frac{2\pi}{3.0} \text{ m}^{-1}$$
$$= 2\pi \times 10^8 / \text{s} = 6.3 \times 10^8 \text{ Hz}$$

$$\omega = 6.3 \times 10^8 \text{ Hz}$$

Finally $\vec{E}_0 = E_0 \hat{i} = 20 \text{ N/C} \hat{i} \Rightarrow \vec{E}_0 = 20 \text{ N/C} \hat{i}$

These completely specify \vec{E} .

$$\begin{aligned} \text{ii) } \vec{\nabla} \times \vec{E} &= \vec{\nabla} \times (\vec{E}_0 e^{i(kz - \omega t)}) \\ &= \vec{\nabla} \times [\vec{E}_0 e^{i(kz - \omega t)}] \end{aligned}$$

Using $\vec{\nabla} \times (f\vec{A}) = f(\vec{\nabla} \times \vec{A}) - \vec{A} \times \vec{\nabla} f$ with $\vec{A} = \vec{E}_0$ gives:

$$\begin{aligned} \vec{\nabla} \times \vec{E} &= -\vec{E}_0 \times \vec{\nabla} (e^{i(kz - \omega t)}) \\ &= -\vec{E}_0 \times \left[\frac{\partial}{\partial x} (\dots) \hat{i} + \dots + \frac{\partial}{\partial z} (\dots) \hat{k} \right] \\ &= -\vec{E}_0 \times ik e^{i(kz - \omega t)} \hat{k} \end{aligned}$$

$$\Rightarrow \vec{\nabla} \times \vec{E} = -ik \left[\underbrace{E_0 \hat{i} \times \hat{k}}_{-\hat{j}} \right] e^{i(kz - \omega t)}$$

$$= -ik E_0 e^{i(kz - \omega t)} \hat{j} = - \frac{\partial \vec{B}}{\partial t}$$

Integrating gives:

$$\vec{B} = \frac{k}{\omega} E_0 e^{i(kz - \omega t)} \hat{j}$$

We see that since we know k, ω, E_0 we can completely specify \vec{B} , once \vec{E} has been specified. Note that

$$k/\omega = 1/c$$

means

$$\vec{B} = \frac{E_0}{c} e^{i(kz - \omega t)} \hat{j}$$

and the amplitude of this is:

$$E_0/c = \frac{20 \text{ N/C}}{3.0 \times 10^8 \text{ m/s}} = 6.7 \times 10^{-8} \text{ T}$$

$$\begin{aligned} \text{c) } \vec{k} \times \vec{E} &= k \hat{k} \times E_0 \hat{i} e^{i(kz - \omega t)} \\ &= k E_0 \hat{j} e^{i(kz - \omega t)} \\ &= \frac{\omega}{c} E_0 e^{i(kz - \omega t)} \hat{j} \\ &= \omega \vec{B} \end{aligned}$$

This illustrates the fact that

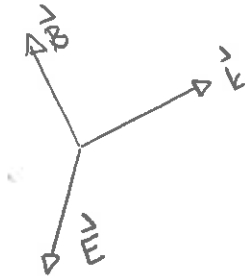
Specifying the electric field completely, results in a complete specification of the magnetic field

via

$$\omega \vec{B} = \vec{k} \times \vec{E}$$

We see that the

Direction of propagation, the electric field and the magnetic field are all perpendicular.



These facts are true for any electromagnetic fields of the form:

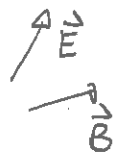
$$\vec{E} = \vec{E}_0 f(\vec{k} \cdot \vec{r} - \omega t)$$

Energy + electromagnetic waves

The energy required to assemble a collection of source charges + source currents can be computed by applying Newtonian mechanics to classical electromagnetism. The resulting energy can be determined from the fields produced by these charges. The result is that:



sources



The energy required to assemble these is:

$$U = \frac{1}{2} \int [\epsilon_0 \vec{E} \cdot \vec{E} + \frac{1}{\mu_0} \vec{B} \cdot \vec{B}] d\tau$$

all space.

We then broaden this to state:

The energy contained within any region of space is:

$$U = \frac{1}{2} \int_{\text{region}} [\epsilon_0 \vec{E} \cdot \vec{E} + \frac{1}{\mu_0} \vec{B} \cdot \vec{B}] d\tau.$$

and the energy per unit volume is

$$u = \frac{\epsilon_0}{2} \vec{E} \cdot \vec{E} + \frac{1}{2\mu_0} \vec{B} \cdot \vec{B}$$

With waves it should be clear that energy can be transported. Specifically with transverse electromagnetic waves we know that

$$\omega B = kE \Rightarrow B = \frac{k}{\omega} E = \frac{E}{c}.$$

for the magnitudes.

Thus the energy density is:

$$u = \frac{\epsilon_0}{2} E^2 + \frac{1}{2\mu_0} B^2 = \frac{\epsilon_0}{2} E^2 + \frac{1}{2\mu_0} \frac{E^2}{c^2}$$

But $c^2 = \frac{1}{\mu_0 \epsilon_0}$ implies:

For transverse electromagnetic waves the energy density is:

$$u = \epsilon_0 E^2 = \frac{B^2}{\mu_0}$$

Again we only need electric field information to describe energy content.

Exercise 2: a) Consider the electromagnetic wave described by

$$\vec{E} = \vec{E}_0 \cos(\vec{k} \cdot \vec{r} - \omega t)$$

where $\vec{k} = k\hat{i}$ and $\vec{E}_0 = E_0\hat{j}$.

Determine an expression for the energy density. Show that this depends on position and time.

b) Suppose

$$\vec{E} = \vec{E}_0 e^{-u^2/2}$$

where $u = \vec{k} \cdot \vec{r} - \omega t$ and $\vec{k} = k\hat{i}$, $\vec{E}_0 = E_0\hat{j}$.

Sketch the energy density as a function of x at $t=0$ s

How does it appear later?

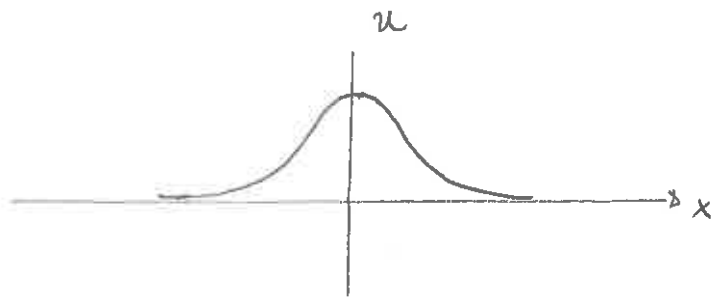
Answer: a) $E^2 = \vec{E} \cdot \vec{E}$
 $= E_0^2 \cos^2(kx - \omega t)$

$\Rightarrow u = \epsilon_0 E_0^2 \cos^2(kx - \omega t)$

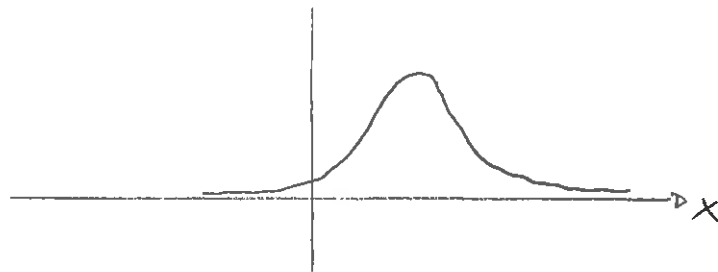
depends on position + time

b) $E^2 = E_0^2 e^{-u^2}$
 $= E_0^2 e^{-(kx - \omega t)^2}$

At $t=0$



Later t



Energy is clearly being transported.