

Mon: HW due

Weds: Read Ch 2.

Damped Oscillations

The solution to the equation of motion for a simple harmonic oscillator is

$$x(t) = C \cos(\omega t + \phi)$$

This describes oscillations that persist indefinitely. The energy of this oscillator is given by

$$E = \alpha \left[\left(\frac{dx}{dt} \right)^2 + \omega^2 x^2 \right]$$

where α is a constant. Substitution from the general solution yields.

$$E = \alpha C^2 \omega^2$$

Exercise: Prove the relationship above. \square

The energy of the oscillator stays constant.

In general observations suggest that oscillations do not persist indefinitely and that energy dissipates. So the equation of motion for a simple harmonic oscillator does not reflect this dissipation. We need to modify it accordingly, so as to yield oscillations that are damped.

Equation of motion for damped oscillations.

Again consider a spring + mass system but suppose that there is an additional drag force proportional to velocity.

We assume that this has horizontal component

$$F_{\text{drag}} = -bv = -b \frac{dx}{dt}$$

where $b > 0$ is a constant:

Newton's 2nd Law gives:

$$\vec{F}_{\text{net}} = m\vec{a}$$

$$\Rightarrow m \frac{d^2x}{dt^2} = -kx - b \frac{dx}{dt}$$

and this gives the equation of motion for a damped oscillator:

$$\boxed{\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = 0}$$

where $\gamma = b/m$ is the damping constant.

$\omega_0 = \sqrt{k/m}$ is the natural frequency.

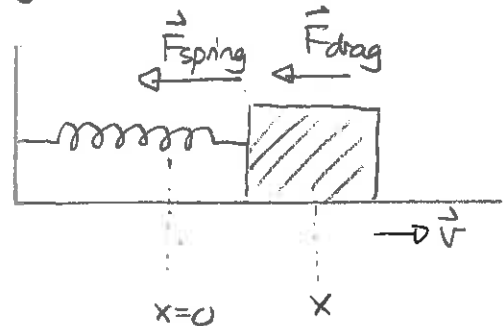
This second order differential equation is again linear.

Solving the equation of motion.

We seek oscillating damped solutions. These have the form

$$x(t) = C e^{-\alpha t} \cos(\omega t + \phi)$$

where $\omega > 0$ is not necessarily ω_0 and $\alpha > 0$ is a constant.



By direct substitution we can show:

1) such solutions only exist when $\gamma < 2\omega_0$

2) $\alpha = \gamma/2$

3) $\omega = \sqrt{\omega_0^2 - \gamma^2/4}$

So we get

If $\gamma < 2\omega_0$ then the general solution to the equation of motion

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = 0$$

is

$$x(t) = C e^{-\gamma t/2} \cos(\omega t + \phi)$$

where

$$\omega = \sqrt{\omega_0^2 - \gamma^2/4}$$

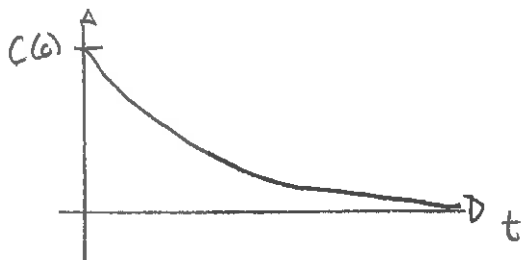
This is called lightly damped or underdamped oscillations. We can rewrite this with a time dependent amplitude.

$$x(t) = C(t) \cos(\omega t + \phi)$$

where

$$C(t) = C e^{-\gamma t/2}$$

The amplitude decays exponentially.



Properties of this solution.

- 1) Period The angular frequency ω differs from the natural frequency and $\omega < \omega_0$.

The previous definition of period, i.e. the shortest time T s.t. $x(t+T) = x(t)$ for all t has no meaning here since the function does not repeat. But we could define the period as

The period is the shortest time between successive moments when the oscillator is at equilibrium and is moving in the same direction.

These occur when

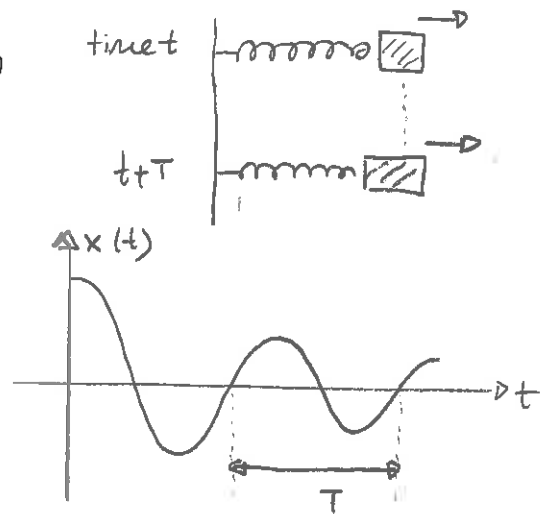
$$x(t) = 0 \Rightarrow \cos(\omega t + \phi) = 0$$

and $\dot{x}(t)$ has same sign. We see that this means

$$\omega(t+T) + \phi = \omega t + \phi + 2\pi$$

$$\Rightarrow \omega T = 2\pi$$

$$\Rightarrow T = \frac{2\pi}{\omega}$$



- 2) Maxima: Maxima occur when

$$\frac{dx}{dt} = 0 \Leftrightarrow C(-\gamma/2)e^{-\gamma t/2} \cos(\omega t + \phi)$$

$$- C\omega e^{-\gamma t/2} \sin(\omega t + \phi) = 0$$

$$\Leftrightarrow \gamma/2 \cos(\omega t + \phi) + \omega \sin(\omega t + \phi) = 0$$

If one maximum occurs at time t' , then at $t'+T$

$$\gamma/2 \cos(\omega(t'+T) + \phi) + \omega \sin(\omega(t'+T) + \phi)$$

$$= \gamma/2 \cos(\omega t' + 2\pi + \phi) + \omega \sin(\omega t' + \phi + 2\pi)$$

$$= \gamma/2 \cos(\omega t' + \phi) + \omega \sin(\omega t' + \phi)$$

But if a max occurs at t' then this expression is zero. So

$$\frac{1}{2} \cos(\omega(t'+\tau) + \phi) + \omega \sin(\omega(t'+\tau) + \phi) = 0$$

which is also a max. Thus

Successive maxima are separated by intervals equal to the period.

Intermediate Dynamics: Group Exercises 6

Damped Harmonic Oscillator

1 Damped harmonic oscillator

The equation of motion for a damped harmonic oscillator is

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \omega_0^2 x = 0. \quad (1)$$

We shall aim to find damped oscillations described by

$$x(t) = Ce^{-\alpha t} \cos(\omega t)$$

where $\alpha > 0$ and $\omega > 0$ are constants (*note that the symbols ω and ω_0 are intentionally different*).

- Assuming this type of solution determine an expression for $\frac{dx}{dt}$.
- Determine an expression for $\frac{d^2x}{dt^2}$.
- Substitute these into Eq. (1) and gather all terms with $\cos(\omega t)$ and all with $\sin(\omega t)$. Determine conditions on α and ω such that the resulting equation is satisfied at all times.
- Is this solution valid for all values of γ and ω_0 ?

Exercise 1

$$\begin{aligned} \text{a) } \frac{dx}{dt} &= C(-\alpha)e^{-\alpha t} \cos(\omega t) + Ce^{-\alpha t}(-\omega) \sin \omega t \\ &= -Ce^{-\alpha t} [\alpha \cos(\omega t) + \omega \sin(\omega t)] \end{aligned}$$

$$\begin{aligned} \text{b) } \frac{d^2x}{dt^2} &= -C(-\alpha)e^{-\alpha t} [\alpha \cos(\omega t) + \omega \sin(\omega t)] \\ &\quad - \neq Ce^{-\alpha t} [-\alpha\omega \sin(\omega t) + \omega^2 \cos(\omega t)] \\ &= Ce^{-\alpha t} \left\{ \cos(\omega t) [\alpha^2 - \omega^2] + \sin(\omega t) 2\alpha\omega \right\} \end{aligned}$$

$$\begin{aligned} \text{c) } & Ce^{-\alpha t} \left\{ [\alpha^2 - \omega^2] \cos(\omega t) + 2\alpha\omega \sin(\omega t) \right\} \\ &+ \gamma(-C)e^{-\alpha t} [\alpha \cos(\omega t) + \omega \sin(\omega t)] \\ &+ \omega_0^2 Ce^{-\alpha t} \cos(\omega t) = 0 \\ \Rightarrow & Ce^{-\alpha t} \left\{ \cos(\omega t) [\alpha^2 - \omega^2 - \gamma\alpha + \omega_0^2] + \sin(\omega t) [2\alpha\omega - \gamma\omega] \right\} = 0 \\ \Rightarrow & \cos(\omega t) [\alpha^2 - \omega^2 - \gamma\alpha + \omega_0^2] + \omega(2\alpha - \gamma) \sin(\omega t) = 0. \end{aligned}$$

At $t=0$ we get $\alpha^2 - \omega^2 - \gamma\alpha + \omega_0^2 = 0$ and this is true for all times.

So $\omega(2\alpha - \gamma) \sin(\omega t) = 0$. This requires $\omega(2\alpha - \gamma) = 0$. So

$$2\alpha - \gamma = 0 \Rightarrow \boxed{\alpha = \gamma/2}$$

$$\begin{aligned} \alpha^2 - \omega^2 - \gamma\alpha + \omega_0^2 = 0 \Rightarrow \gamma^2/4 - \omega^2 - \gamma^2/2 + \omega_0^2 = 0 \Rightarrow \omega^2 &= \omega_0^2 - \gamma^2/4 \\ &= \sqrt{\omega_0^2 - \gamma^2/4}. \end{aligned}$$

$$\text{d) No we need } \omega_0^2 \geq \gamma^2/4 \Rightarrow \boxed{\omega_0 > \gamma/2}$$

2 Lightly damped oscillations

The amplitude for a lightly damped oscillator is

$$C(t) = Ce^{-\gamma t/2}$$

where C is a constant.

- Determine $C(1\text{ s})/C(0\text{ s})$.
- Determine $C(2\text{ s})/C(1\text{ s})$.
- Determine $C(t + 1\text{ s})/C(t)$ for any time t .
- Determine γ from the ratio $C(t + 1\text{ s})/C(t)$ for any time t .
- Determine γ from the ratio $C(t + T)/C(t)$ where T for any time t but where T is the period of oscillation.

Answer:

a) *From the definition*

$$\frac{C(1\text{ s})}{C(0\text{ s})} = \frac{C_0 e^{-\gamma/2}}{C_0 e^{-\gamma \cdot 0/2}} = e^{-\gamma/2}$$

b) *From the definition*

$$\frac{C(2\text{ s})}{C(1\text{ s})} = \frac{C_0 e^{-\gamma \cdot 2/2}}{C_0 e^{-\gamma/2}} = \frac{e^{-\gamma}}{e^{-\gamma/2}} = e^{-\gamma + \gamma/2} = e^{-\gamma/2}$$

c) *From the definition*

$$\frac{C(t + 1\text{ s})}{C(t)} = \frac{C_0 e^{-\gamma(t+1)/2}}{C_0 e^{-\gamma t/2}} = \frac{e^{-\gamma t/2 - \gamma/2}}{e^{-\gamma t/2}} = e^{-\gamma t/2 - \gamma/2 + \gamma t/2} = e^{-\gamma/2}$$

Notice that the fractional reduction in amplitude over any 1 s interval is the same regardless of when the interval started. Then

$$\ln\left(\frac{C(t + 1\text{ s})}{C(t)}\right) = \ln(e^{-\gamma/2})$$

gives

$$\ln\left(\frac{C(t + 1\text{ s})}{C(t)}\right) = -\frac{\gamma}{2}$$

and thus

$$\gamma = -2 \ln\left(\frac{C(t + 1\text{ s})}{C(t)}\right) = 2 \ln\left(\frac{C(t)}{C(t + 1\text{ s})}\right)$$

d) A similar derivation gives

$$\frac{C(t+T)}{C(t)} = \frac{C_0 e^{-\gamma(t+T)/2}}{C_0 e^{-\gamma t/2}} = e^{-\gamma T/2}$$

Again, using natural logarithms,

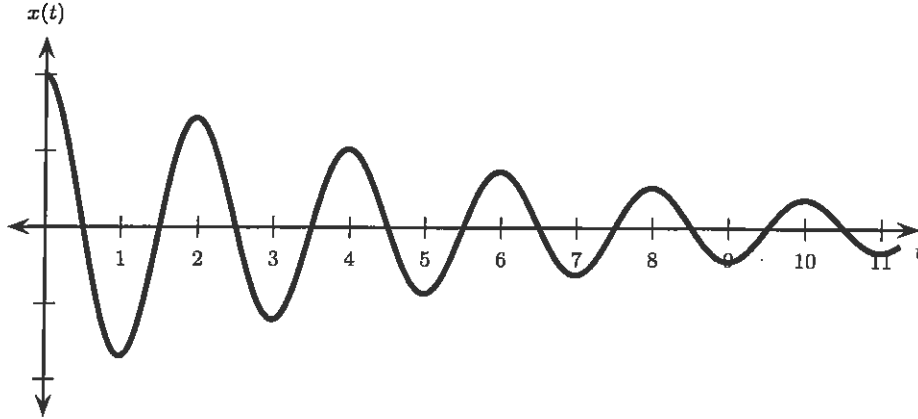
$$-\frac{\gamma T}{2} = \ln\left(\frac{C(t+T)}{C(t)}\right)$$

and this gives

$$\gamma = -\frac{2}{T} \ln\left(\frac{C(t+T)}{C(t)}\right) = \frac{2}{T} \ln\left(\frac{C(t)}{C(t+T)}\right)$$

3 Determining the damping constant

A graph of position vs. time for a lightly damped oscillator is shown. The time is measured in seconds.



Determine γ using the data from the graph.

Answer:

In general the amplitude satisfies

$$C(t) = C_0 e^{-\gamma t/2}$$

Then

$$\gamma = -\frac{2}{T} \ln \left(\frac{C(t+T)}{C(t)} \right)$$

Here, with $t = 0$ s and $T = 4$ s, inspection from the graph gives

$$\frac{C(t+T)}{C(t)} = \frac{1}{2}$$

Thus

$$\gamma = -\frac{2}{T} \ln \left(\frac{C(t+T)}{C(t)} \right) = -\frac{2}{4 \text{ s}} \ln \left(\frac{1}{2} \right) = 0.35 \text{ s}^{-1}$$