# Correlated quantum states and enhanced mixed state Pauli channel parameter estimation



### Abstract

The accuracy of any physical scheme used to estimate parameters that govern the evolution of quantum systems is limited by statistical fluctuations inherent in quantum measurement processes. Quantum estimation theory provides methods for quantitative comparison of various estimation protocols. We focus on estimating the parameter governing the single qubit Pauli channel, for which it is known that optimal pure state estimation uses unentangled initial states. We consider a restricted version of this problem in which the initial states of the individual quantum systems are not pure and ask whether there are quantum estimation protocols which can yield greater accuracies than independent channel use protocols, analogous to classical repetition and averaging schemes. We compare a protocol involving quantum correlated states to independent channel use protocols. We show, that unlike the pure state case, the quantum correlated state protocol can yield greater estimation accuracy than any independent state protocol. We show that these gains persist even when the system states are separable and, in some cases, when quantum discord is absent after channel invocation.

The evolution of any quantum system is described by a quantum operation, which may depend on one or more parameters. For example, the **Pauli channel** maps a single qubit as

$$\hat{\rho} \stackrel{\Gamma(\lambda)}{\mapsto} \hat{\rho}_{\mathsf{f}}(\lambda) := (1 - \lambda) \,\hat{\rho} + \lambda \hat{\sigma}_n \hat{\rho} \hat{\sigma}_n$$

and depends on the parameter  $0 \leq \lambda \leq 1$  and  $\hat{\sigma}_n$  is a Pauli operator with n representing a direction.

### Task: Knowing the direction, n, estimate parameter, $\lambda$ , as accurately as possible by subjecting quantum systems to the channel.

Any physical estimation protocol requires preparation of input states for system qubits, followed by channel invocations on some or all of the qubits and terminates in measurements on qubits. The outcomes of measurements can be applied to an estimator function, which returns an estimate,  $\lambda_{est}$ , for  $\lambda$ . The probabilistic nature of outcomes of measurements on quantum systems and the effects of measurements on quantum states imply that repeated runs of any physical estimation protocol will yield estimates which fluctuate around the true parameter value.

## Quantifying Estimation Accuracy: Fisher Information

The accuracy of the measurement is quantified in terms of the mean square error,

m. s. e  $(\lambda_{\text{est}}) = \left\langle \left(\lambda_{\text{est}} - \lambda\right)^2 \right\rangle$ .

For any unbiased estimator, the **Cramér-Rao bound** gives [1]:



Classical Fisher information depends on measurement choice but not estimator.

The quantum Fisher information is given by

$$H(\lambda) = \operatorname{Tr}\left(\hat{\rho}_{\mathsf{f}}\hat{L}^{2}\right) \qquad \text{with} \qquad \frac{\partial\hat{\rho}_{\mathsf{f}}}{\partial\lambda} = \frac{1}{2}\left(\hat{\rho}_{\mathsf{f}}\hat{L} + \hat{L}\hat{\rho}_{\mathsf{f}}\right)$$

**Optimal estimation: choose input state and additional parameter-independent** unitaries so as to maximize the quantum Fisher information.

The quantum Fisher information can be increased by invoking the channel multiple times. If there are m invocations on identically prepared independent or uncorrelated quantum systems, then it is always true that  $H(\lambda) = mH_{\rm s}(\lambda)$  where  $H_{\rm s}(\lambda)$  is the quantum Fisher information for one invocation on one systems.

Quantum estimation: Given m channel invocations is it possible, by using quantum resources, to exceed the *m*-fold increase attained with "classical" independent uses of the channel?

The independent channel use improvement in the QFI can be enhanced by an additional factor of m for certain scenarios by using entangled or correlated states [2, 3].

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correlations cannot enhance parameter estimation accuracy.

is a phase flip channel.

Channel invocation on m qubits.



where  $\hat{
ho}_j = \frac{1}{2} \left( \hat{I} + \mathbf{r}_j \cdot \hat{\boldsymbol{\sigma}} \right)$ . Polarization,  $r := |\mathbf{r}_{i}|$ , is the same for all qubits.

Channel invocation on m of n qubits is preceded by a correlating preparatory unitary.







$$H(\lambda) = \frac{m^2(1-2\lambda)^{2m-2}}{2^{n-1}} \sum_{j=0}^n \binom{n}{j} \frac{1}{c_{j+1}^2} \frac{1}{2^{n-1}} \frac{1}{2^{n-1}}$$

where

$$c_{j\pm} = (1+r)^j (1-r)^{n-j} \pm (1+r)^{j}$$

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